# <u>Kirchoff's Voltage Law</u>

Consider a simple electrical circuit:



We find that if the **voltage source** is on (i.e.,  $V \neq 0$ ), then there will be electric potential differences (i.e., voltage) between different points of the circuit. This can **only** be true if **electric fields** are present!

The electric field in this circuit will "look" something like this:

So, instead of using circuit theory, let's use our new electromagnetic knowledge to analyze this circuit.

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First, consider a **contour**  $C_1$  that follows the circuit path.

 $V_1$ 



Using this path, let's evaluate the contour integral:

$$\oint_{C_1} \mathbf{E}(\overline{r}) \cdot \overline{d\ell}$$

This is most easily done by breaking the contour  $C_1$  into six sections: section 1 extends from point *a* to point *b*, section 2 extends from point *b* to point *c*, etc. Thus, the integral becomes:

$$\oint_{C_1} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} = \int_{a}^{b} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} + \int_{b}^{c} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} + \int_{c}^{d} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} +$$

Let's evaluate each term individually:

#### $\leq$ Section 1 (a to b)

In this section, the contour follows the wire from the voltage source to the first resistor. We know that the electric field in a perfect conductor is zero, and likewise in a good conductor it is very small. Assuming the wire is in fact made of a good conductor (e.g. copper), we can approximate the electric field within the wire (and thus at every point along section 1) as zero (i.e.,  $\mathbf{E}(\vec{r}) = 0$ ). Therefore, this first integral equals zero!

$$\int_{a}^{b} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} = 0$$

This of course makes sense! We know that the electric potential difference across a **wire** is **zero volts**.

#### Section 2 (b to c)

In this section, the contour moves through the first **resistor**. The contour integral along this section therefore allows us to determine the electric **potential difference** across this resistor. Let's denote this potential difference as  $v_1$ :

$$\int_{b}^{c} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} = V(P_{b}) - V(P_{c}) = v_{1}$$

#### Section 3 (c to d)

Just like section 1, the contour follows a **wire**, and thus the electric field along this section of the contour is **zero**, as is the potential difference between point *c* and point *d*.

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$$\int \mathbf{E}(\bar{r})\cdot \overline{d\ell} = \mathbf{0}$$

#### Section 4 (d to e)

Just like section 2, the contour moves through a **resistor**. The contour integral for this section is thus equal to the potential difference across this **second** resistor, which we denote as  $v_2$ :

$$\int_{d} \mathbf{E}(\bar{r}) \cdot \bar{d\ell} = \mathbf{V}(P_{d}) - \mathbf{V}(P_{e}) = \mathbf{v}_{2}$$

## ) \$ <u>Section 5 (e to f</u>)

Again, the contour follows a conducting **wire**—and again, the electric field along the contour and the potential difference across it are both **zero**:

$$\int \mathbf{E}(\overline{r}) \cdot \overline{d\ell} = 0$$

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### Section 6 (f to a)

This **final** section of contour  $C_1$  extends through the **voltage source**, thus the contour integral of this section provides the electric potential difference between the two terminals of the this voltage source (i.e.,  $V(P_f) - V(P_a)$ ). By **definition**, the potential difference between points *a* and *f* is a value of V volts (i.e.,  $V(P_a) - V(P_f) = V$ ). Therefore, we find that the contour integral of section 6 is :

$$\int_{f}^{a} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} = V(P_{f}) - V(P_{a})$$
$$= -(V(P_{a}) - V(P_{f}))$$
$$= -V$$

Whew! Now let's combine these results to determine the contour integral for the entire contour  $C_1$ .

$$\oint_{\mathcal{C}_{1}} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} = \int_{a}^{b} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} + \int_{b}^{c} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} + \int_{c}^{d} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} + \int_{c}^{d} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} + \int_{c}^{d} \mathbf{E}(\bar{r}) \cdot \overline{d\ell} + \int_{c}^{a} \mathbf{E}(\bar{r}) \cdot \overline{$$

**Q:** Wait; I've forgotten, Why are we evaluating these contour integrals ?

A: Remember, since the electric field is static, we also know that integral around any closed contour is zero. Thus, we can conclude that:

$$\mathbf{0} = \oint_{\mathcal{C}_1} \mathbf{E}(\overline{\mathbf{r}}) \cdot \overline{\mathbf{d}\ell} = \mathbf{v}_1 + \mathbf{v}_2 - \mathbf{V}$$

In other words, we find by performing an **electromagnetic** analysis of the circuit, the voltages across each circuit element are related as:

 $v_1 + v_2 - V = 0$ 

**Q:** You have wasted my time! Using only Kirchoff's Voltage Law (KVL), **I** arrived at precisely the same result  $(v_1 + v_2 - V = 0)$ . **I** think the above equation is true because of KVL, not because of your fancy electromagnetic theory!

A: It is true that the result we obtained by integrating the electric field around the circuit contour is likewise apparent from KVL. However, this result is still attributable to electrostatic physics, because KVL is a direct result of electrostatics!

Jim Stiles

The electrostatic equation :

$$\oint_{C} \mathbf{E}(\bar{\mathbf{r}}) \cdot \overline{d\ell} = \mathbf{0}$$

when applied to the closed contour of any **circuit**, results in **Kirchoff's Voltage Law**, i.e.:

$$\sum_{n} v_{n} = 0$$

where  $v_n$  are the electric potential differences across each element of a circuit "loop" (i.e., closed contour).

Gustav Robert Kirchhoff (1824-1887), German physicist, announced the laws that allow calculation of the currents, voltages, and resistances of electrical networks in 1845, when he was only **twenty-one**! His other work established the technique of spectrum analysis that he applied to determine the composition of the Sun.

From www.ee.umd.edu/~taylor/frame5.htm

